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TECHNICAL REPORT BRL-TR-2672

A COMPUTATIONAL INVESTIGATION OF THE EFFECT OF SHIELDING IN MITIGATING SHOCK INITIATION STIMULI PRODUCED BY IMPACT

John Starkenberg Toni M. Dorsey Kelly J. McGlothlin

September 1985



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We have conducted a study of the role of shielding in reducing the shock initiation stimulus for a simple one-dimensional representation of the problem of sympathetic detonation of munitions. We found that single layered shields made of materials with low acoustic impedance generally produce a complex shock wave structure in the acceptor. This complex structure is associated with a low level of initiation stimulus because of the breakup into several					
weaker shocks and the reduction of $\int_{p^2}^{p^2} dt$. High impedance shield materials.					

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ABSTRACT (continued)

the integral OI P squared with respect to + also substantially reduce the initiation stimulus, but without the accompanying shock breakup effect. Increasing shield thickness improves performance and can change the order of effectiveness of shield materials. With multilayered shields composed of a high-impedance and a low-impedance material we observed shock structures depending strongly on the ordering of the materials in the shield. Multiple shock structure was usually observed when the highimpedance material was the outer component of the shield and sufficient lowimpedance material was present. Single compression waves with variable peak pressures and rise times were usually observed when the low-impedance material was the outer component of the shield. Substantial benefits in terms of

shock breakup and p dt reduction can be obtained by increasing the thickness of three-layered shields, which were found to perform better than five-layered shields.

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I. INTRODUCTION

In order to reduce the vulnerability of stored ammunition, shielding between rounds has been used to prevent initiation of detonation in a round (the acceptor) when one of its neighbors (the donor) detonates. The shielding serves to prevent direct impact of the donor casing or its fragments on the acceptor as well as to process the shock wave entering the acceptor, thus reducing the initiation stimulus. The latter mechanism is amenable to analysis using a hydrodynamic computer code and we have already simulated round to round detonation propagation, with and without shielding, using the 2DE code. 1,2 Shielding effectiveness in reducing shock initiation stimulus levels is further amenable to analysis in one dimension. The present report, therefore, concerns our one-dimensional study of shielding effectiveness using the STEALTH code. 3

TI. SIMULATION DESCRIPTION

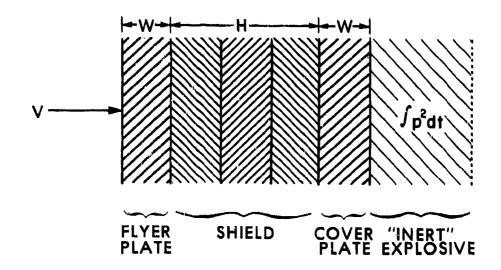
The problem elements are the donor, the shield and the acceptor. The acceptor, in the one-dimensional simulation, consists of a steel layer (cover plate or acceptor casing) covering a layer of "inert" TNT. The shield configuration is varied. It consists of a single layer of any of various materials or multiple layers of two different materials. The donor is represented by a layer of steel (flyer plate or donor casing) with an initial velocity on the order of that achieved by the casing of a detonating round. The flyer and cover plate thicknesses are always equal. The problem geometry is illustrated in Figure 1. The "inert" TNT is described by Lee's unreacted JWL equation of state and all other materials by Wilken's LLNL model with 304 steel used for the flyer and cover plates.

¹P.M. Howe, Y.K. Huang and A.L. Arbuckle, "A Numerical Study of Detonation Propagation Between Munitions," Seventh Symposium (International) on Detonation, pp. 1055-1061, June 1981.

²J. Starkenberg, Y.K. Huang and A.L. Arbuckle, "A Two-Dimensional Numerical Study of Detonation Propagation Between Munitions by Means of Shock Initiation," BRL Technical Report ARBRL-TR-02522, September 1983.

³STEALTH-A Lagrange Explicit Finite Difference Code for Solids, Structural, and Thermohydraulic Analysis, Electric Power Research Institute, November 1981.

⁴E.L. Lee and C.M. Tarver, "Phenomenological Model of Shock Initiation in Heterogeneous Explosives," Physics of Fluids, Volume 23, Number 12, pp. 2362-2372, December 1980.



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Figure 1. One-Dimensional Representation.

III. CHARACTERIZATION OF THE SHOCK INITIATION STIMULUS

The simplest characterization of shock initiation stimulus levels is the square of the pressure integrated with respect to time ($\int p^2 dt$) evaluated in the acceptor explosive. The use of this integral is an extension of the critical energy concept which is applicable to singly shocked explosives in planar experiments. We have, in general, used this parameter to characterize our computed results. However, the shields often alter the rise time of a single compression wave or produce multiple shock loading in the acceptor. In these cases, the ramp wave or shock waves after the first are not presumed to contribute as significantly as a single shock to the initiation. Thus, integrating through the ramp or including all shocks in the $\int p^2 dt$ calculation produces a conservative estimate of the stimulus (i.e. the actual conditions are even more predisposed toward preventing acceptor initiation). It is also instructive to consider in detail the processing of the shock wave produced by the shield.

IV. RESPONSE OF UNSHIELDED ACCEPTORS

In order to provide baseline data from which to evaluate the stimulus reduction provided by shielding, we ran a number of computations without shielding. This type of loading always produces a single shock in the acceptor as illustrated in Figure 2. In these we varied the casing thickness and impact velocity so as to produce sets of results for constant velocity, constant momentum and constant energy impact. Casing thickness was varied from 5 to 20 mm with appropriate velocities. The results are summarized in Table 1 and plotted in Figure 3. Not surprisingly, as the casing thickness is increased at constant velocity, the initiation stimulus increases rapidly. For constant momentum impacts the stimulus decreases rapidly with increasing casing thickness. Constant energy impacts do not produce a constant stimulus, rather the initiation stimulus decreases slowly with increasing casing thickness. The results of these computations suggest a correlation with $\frac{3}{5}$ V, as illustrated in Figure 4.

Table 1. Initiation Stimulus With Unshielded Acceptors

	CONSTANT	VELOCITY	CONSTANT N	MOMENT UM	CONSTANT	ENERGY
CASING THICKNESS W (mm)	VELOCITY V (km/s)	$\int_{\rho}^{2} dt$ (GPa ² -ms)	VELOCITY V (km/s)	$\int_{p^2 dt}^{2} dt$ $GPa^2 - ms$	VELOCITY V (km/s)	$\int_{p^2 dt}^{2} dt$ (GPa ² -ms)
5.0	1.00	0.074	-	-	1.41	0.170
7.5	1.00	0.112	1.33	0.222	1.15	0.157
10.0	1.00	0.145	1.00	0.145	1.00	0.145
12.5	1.00	0.171	0.80	0.100	0.89	0.129
15.0	1.00	0.190	0.67	0.074	0.82	0.119
17.5	1.00	0.208	0.57	0.054	0.76	0.108
20.0	1.00	0.222	0.50	0.042	0.71	0.098

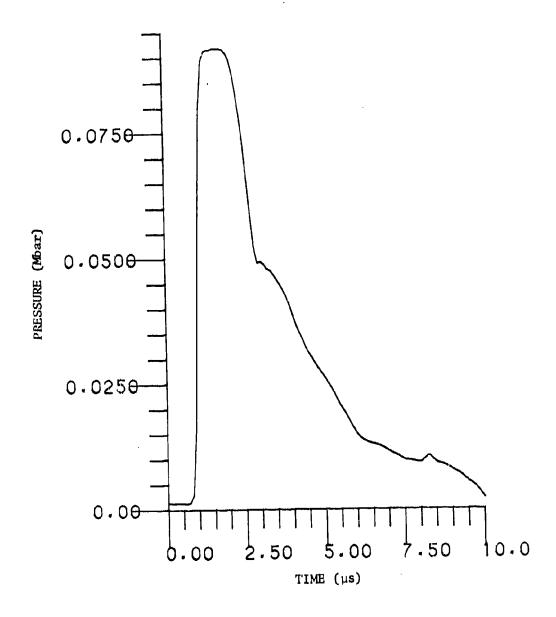


Figure 2. Typical Acceptor Explosive Shock Pressure History without Shielding

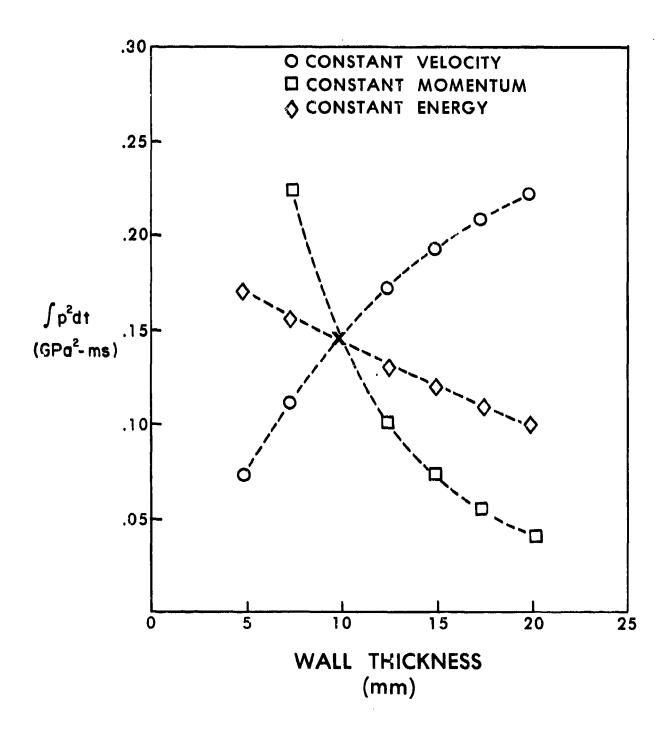


Figure 3. Variation of Shock Initiation Stimulus with Wall Thickness in the Absence of Shielding

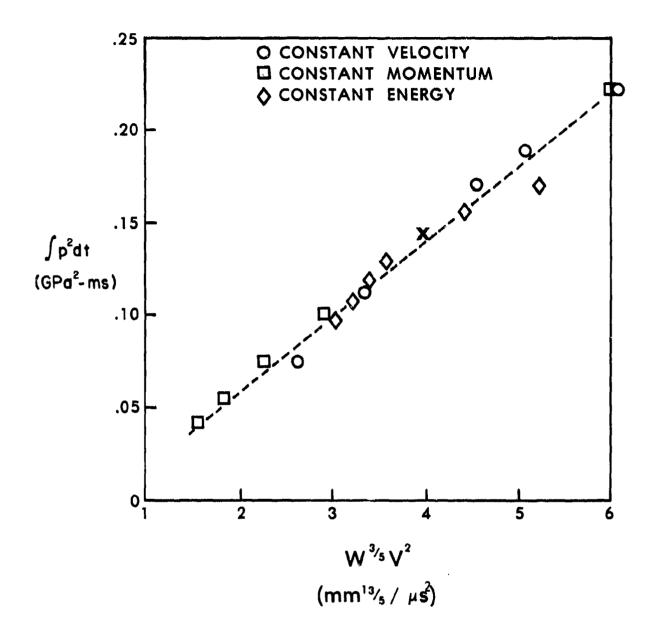


Figure 4. Correlation of Shock Initiation Stimulus with Flyer and Cover Plate Parameters.

V. RESULTS WITH SINGLE LAYERED SHIELDS

A. General

Our version of the STEALTH code is accompanied by a materials library providing standard preprogrammed material models. The vast majority of these are for metallic elements and alloys. The only plastic material description available is for Lucite. Nonetheless, we made computations with 10 mm thick casings and various shield thicknesses for each available material. Special emphasis was given to Lucite, tungsten and steel. The impact velocity was held at 1 km/s.

B. Shock Structure

The pressure history in the acceptor explosive depends on the casing thickness, the shield thickness and the shock impedance of the shield. When the impedance of the shield is lower than that of the casing, a shock wave reverberates between the donor and acceptor casings, delivering repeated shocks to the acceptor. Our lowest impedance material is Lucite. Figure 5 illustrates this effect for 5, 15, 25 and 35 mm thick Lucite shields. Clear separation of three shocks of increasing amplitude may be observed in Figure 5c for a 25 mm thick shield. The complete pulse is about 14 μs long with a peak pressure of 35 GPa (.035 Mbar). For the 15 mm thick shield of Figure 5b, the first shock is clearly defined while the second shock appears intermingled with the third. This pulse is about 12 μs in duration with a 4.2 GPa (.042 Mbar) peak. With a 5 mm thick shield, all the shocks have coalesced into a 10 ms pulse with a 5.1 GPa (.051 Mbar) peak, as illustrated in Figure 5a. Thus the effect of increasing shield thickness is to increase the interval between shocks, lengthen the pulse duration and reduce the peak pressure. The temporal spacing between the shock fronts generated at the shield/acceptor casing interface depends on shield thickness since two shock transits across the shield occur between the generating interactions. As these shocks propagate into the acceptor, they tend to coalesce so that, if the initial temporal spacing is small enough, the shocks will not be distinguishable from one another when they arrive at the acceptor explosive. Because of this coalescence, shock breakup was only observed for Lucite and magnesium (see Figure 6).

When the impedance of the shield is higher than that of the casing a single shock in the acceptor is always observed. Our highest impedance material is tungsten. Figure 7 shows the shocks produced with 5, 15, 25 and 35 mm thick tungsten shields. The effect of increasing shield thickness is to slightly increase pulse duration while substantially reducing peak pressure.

C. Initiation Stimulus Reduction

The effect of shield material on shock initiation stimulus reduction as measured by $\int p^2 dt$ is illustrated in Figure 8. This is a plot of $\int p^2 dt$ versus initial acoustic impedance for 30 mm thick shields. The results are segregated into two groups. One group includes most of the materials whose initial acoustic impedance is less than that of steel and the other group is comprised primarily of materials whose initial acoustic impedance is greater than that of steel. Exceptions are that nickel with a slightly greater

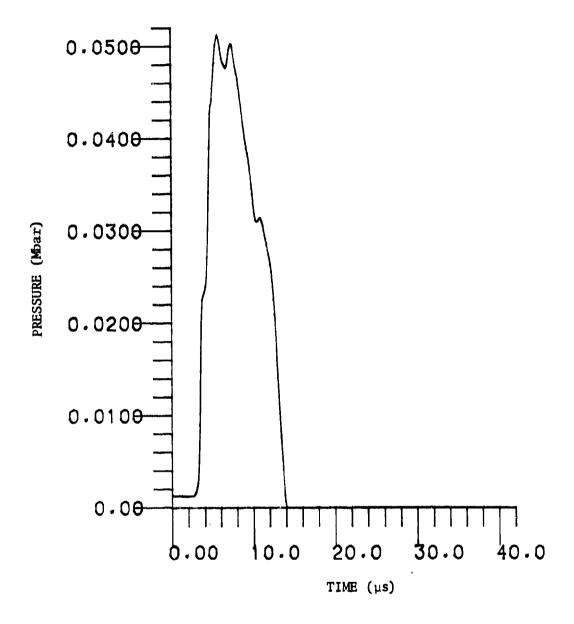


Figure 5a. Acceptor Explosive Pressure History with 5-mm Thick Lucite Shield

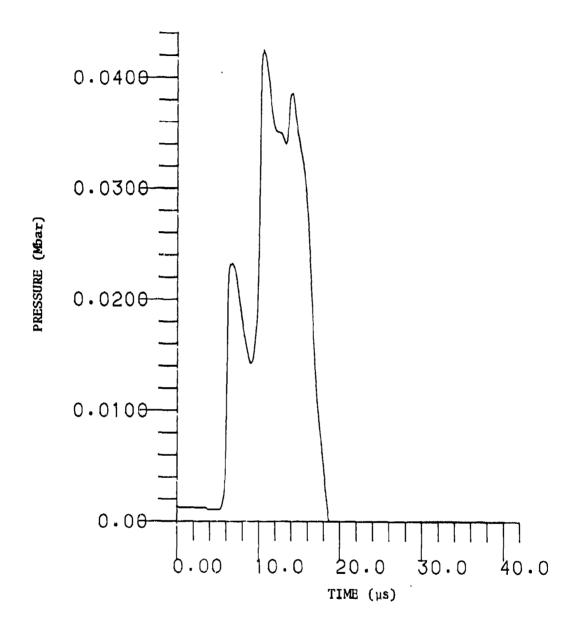


Figure 5b. Acceptor Explosive Pressure History with 15-mm Thick Lucite Shield

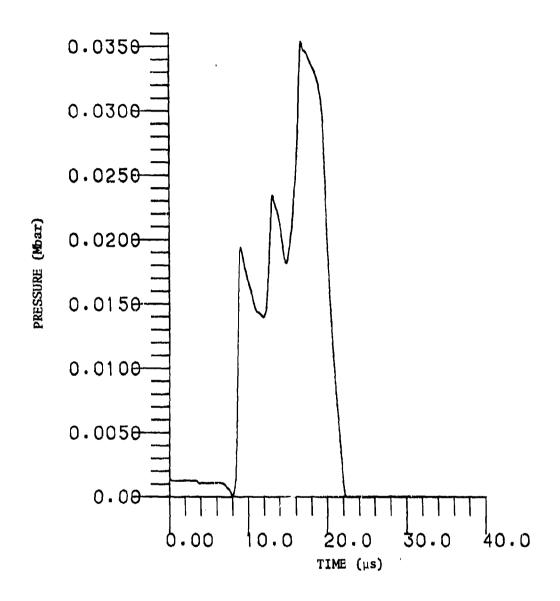


Figure 5a. Acceptor Explosive Pressure History with 25-mm Thick Lucite Shield

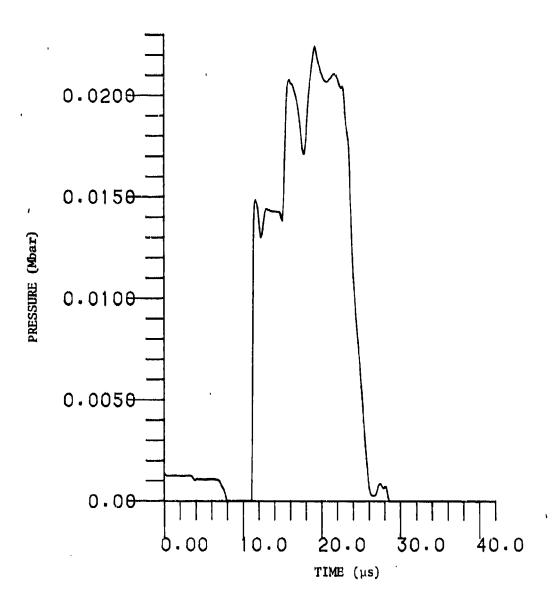


Figure 5d. Acceptor Explosive Pressure History with 35-mm Thick Lucite Shield

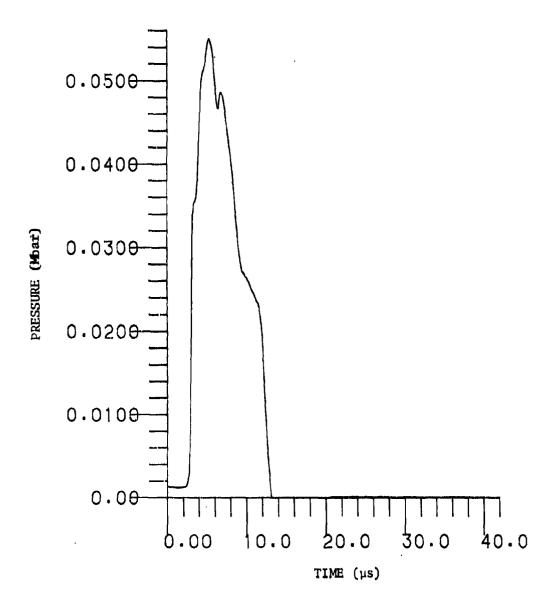


Figure 6a. Acceptor Explosive Pressure History with 5-mm Thick Magnesium Shield

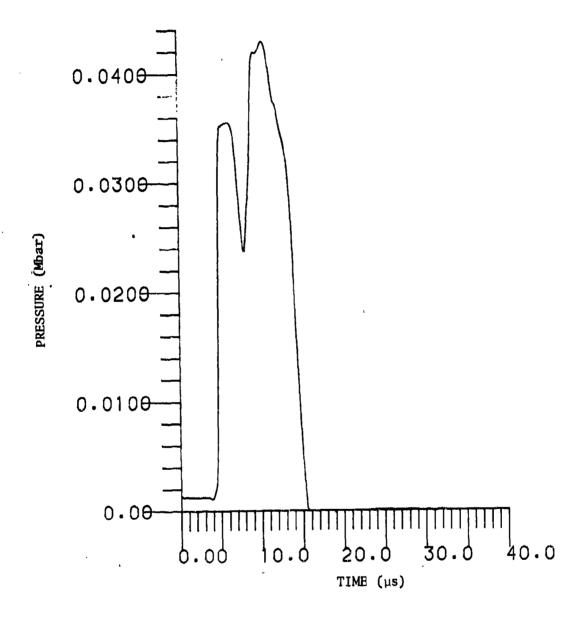


Figure 6b. Acceptor Explosive Pressure History with 15-mm Thick Magnesium Shield

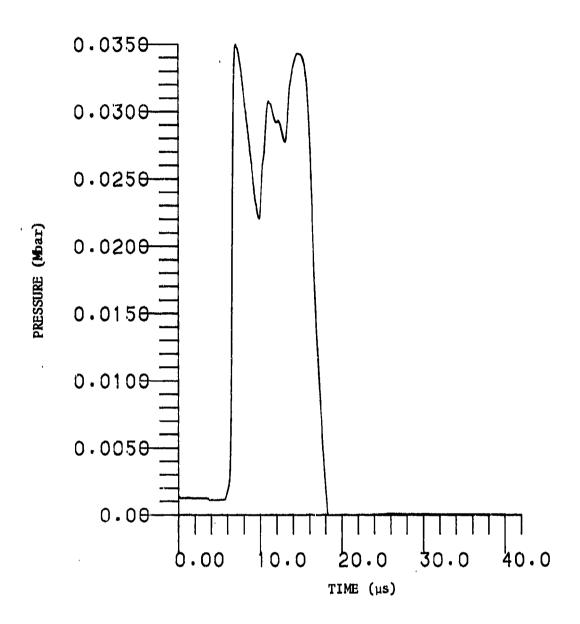


Figure 6c. Acceptor Explosive Pressure History with 25-mm Thick Magnesium Shield

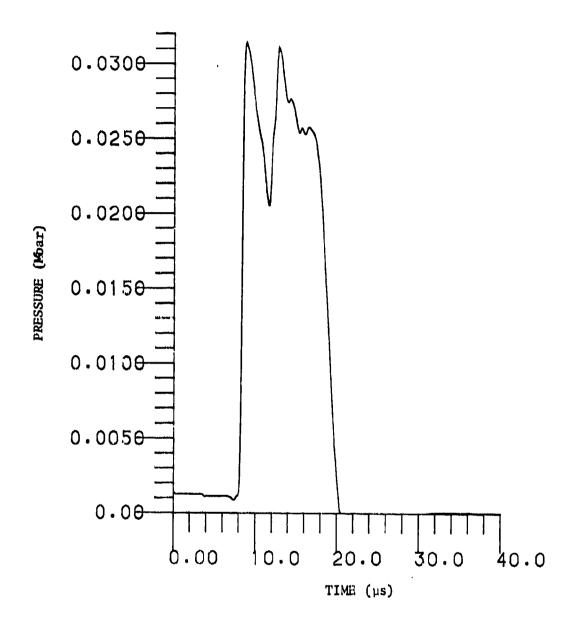


Figure 6d. Acceptor Explosive Pressure History with 35-mm Thick Magnesium Shield

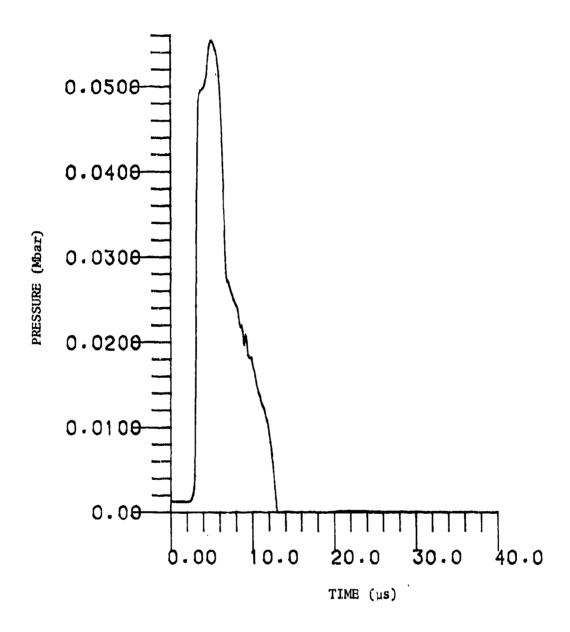


Figure 7a. Acceptor Explosive Pressure History with 5-mm Thick Tungsten Shield

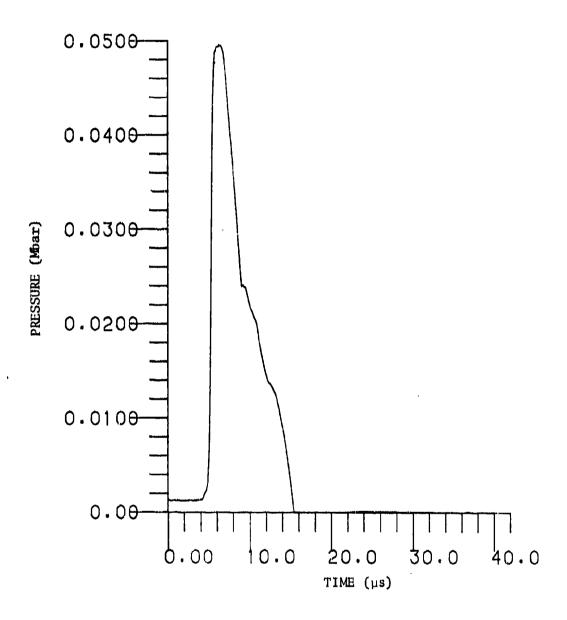


Figure 7b. Acceptor Explosive Pressure History with 15-mm Thick Tungsten Shield

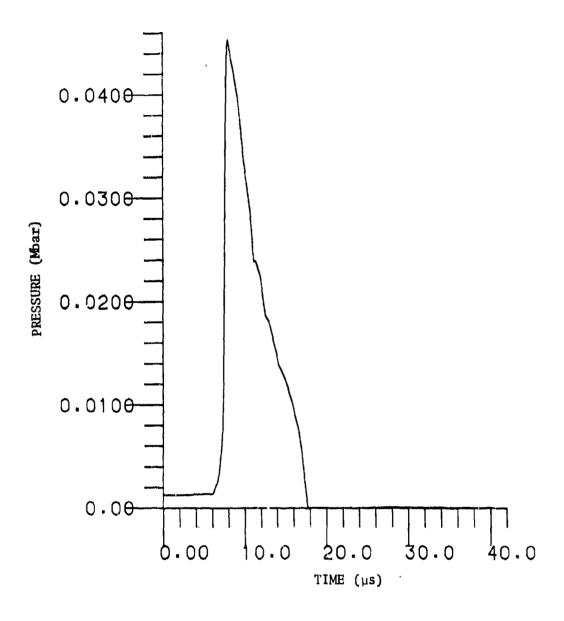


Figure 7c. Acceptor Explosive Pressure History with 25-mm Thick Tungsten Shield

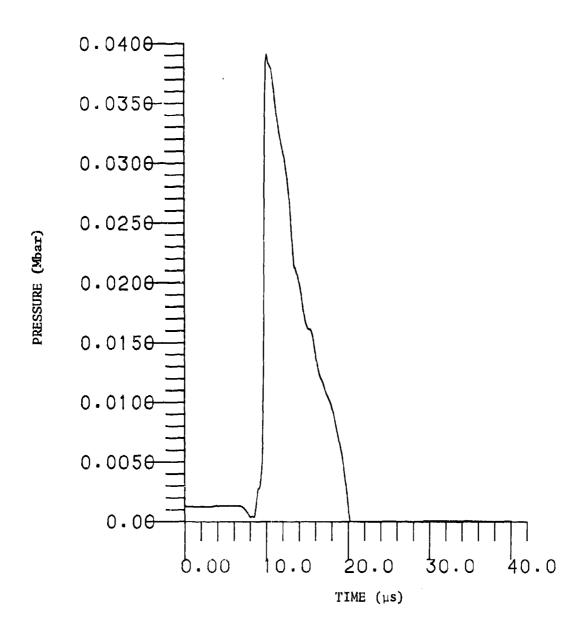


Figure 7d. Acceptor Explosive Pressure History with 35-mm Thick Tungsten Shield

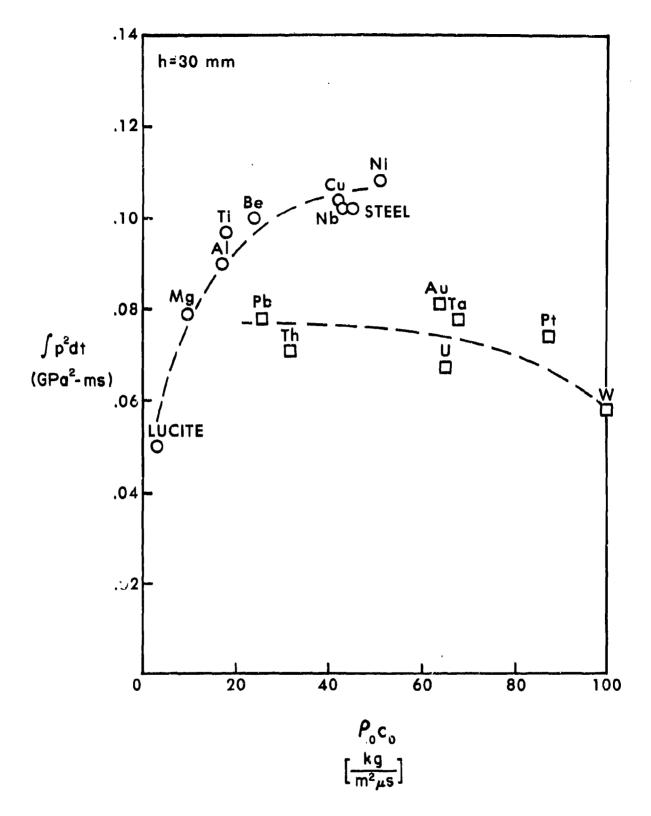


Figure 8. Variation of Shock Initiation Stimulus with Initial Acoustic Impedance of 30-mm Thick Shields

initial impedance than steel belongs to the low impedance group and lead and thorium with lower impedance than steel appear to belong to the higher impedance group (although the latter classification is more questionable). As impedance increases among the low impedance materials, the initiation stimulus approaches that associated with steel shields, which provide the lowest level of protection. Materials in the high impedance group, on the other hand, do not appear to produce stimulus levels approaching that of steel, but they provide greater protection than all but the lowest impedance materials.

The effect of shield thickness is illustrated in Figure 9 which is a plot of $\int p^2 dt$ versus shield thickness for Lucite, steel and tungsten. The results show that steel provides the least protection, except for shields less than about 6 mm thick where Lucite is worst, and tungsten provides the best protection except for shields between about 27 mm and 50 mm thick where Lucite is slightly better. Stimulus reduction increases with shield thickness for all three materials. The increase is most marked for Lucite, which also shows a leveling off for shield thicknesses greater than 40 mm at which point little additional protection is provided with increasing thickness.

VI. RESULTS WITH MULTILAYERED SHIELDS

A. General

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Shields with multiple layers provide additional impedance discontinuities which can further reduce the shock initiation stimulus. We have examined the effectiveness of three and five layer shields composed of alternate layers of our lowest and highest impedance materials, Lucite and tungsten, and of alternate layers of Lucite and Steel as well. The order of the materials and the relative thickness of the layers was varied but, except in the study of shield thickness, all of the shields were 30 mm thick. Symmetry was always maintained. The impact velocity was held at 1 km/s.

When the total shield thickness is fixed, three-layered shields may be completely specified by the Lucite thickness fraction (i.e. the total thickness of all Lucite layers divided by the shield thickness, h) with five-layered shields, an additional degree of freedom arises since the thickness of the innermost layer need not equal the thickness of the two outer layers of the same material. This degree of freedom is accounted for by the parameter $z=h_1/(h-h_0)$ where h_1 is the thickness of the central layer and h_0 the thickness of the outer layers. When z=0, the shield is reduced to three layers; and when $z=h_1/h$, the shield also has three layers but with the material order reversed.

B. Shock Structure

Three-layered shields composed of Lucite and tungsten may have either component in the inner layer. Figure 10 illustrates the pressure history in the acceptor for shields consisting of an inner layer of Lucite surrounded by layers of tungsten. When the shield is composed mostly of tungsten, as in Figure 10a, a single shock enters the acceptor. Figure 10b shows results for a shield with a little more Lucite. Some structure at the tail of the wave

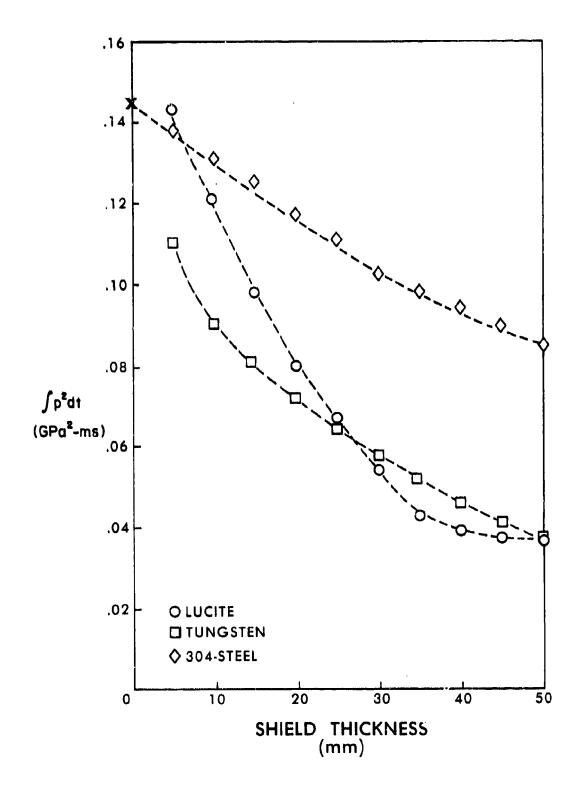


Figure 9. Variation of Shock Initiation Stimulus with Shield Thickness for Lucite, Tungsten and 304-Steel Shields

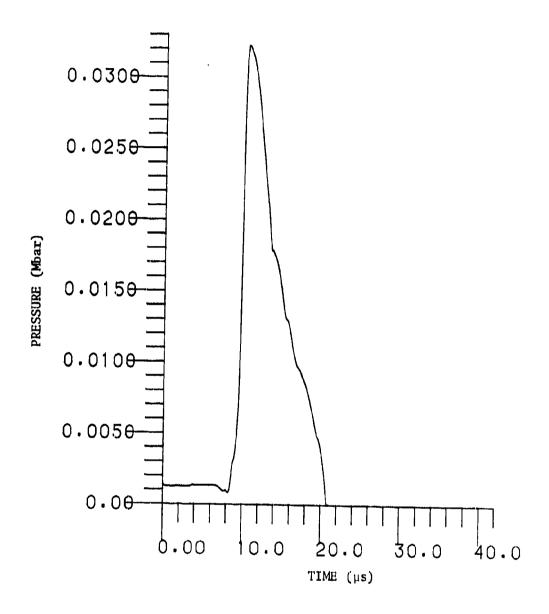


Figure 10a. Acceptor Explosive Pressure History with 30-mm Thick Three-Layered Tungsten/Lucite/Tungsten Shield at 8.3% Lucite

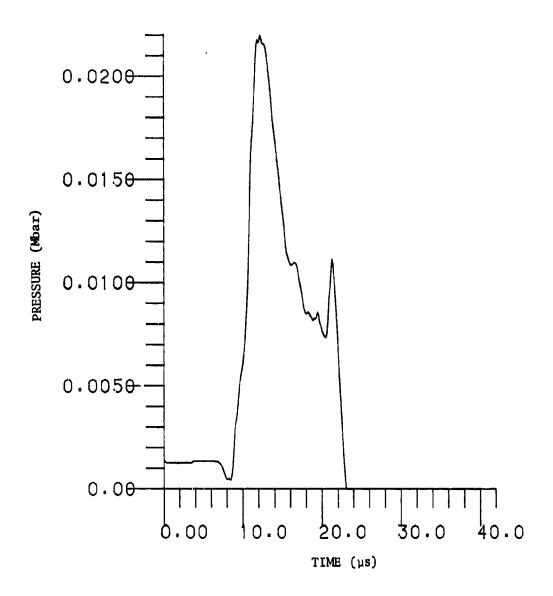


Figure 10b. Acceptor Explosive Pressure History with 30-mm Thick
Three-Layered Tungsten/Lucite/Tungsten Shield at 16.7% Lucite

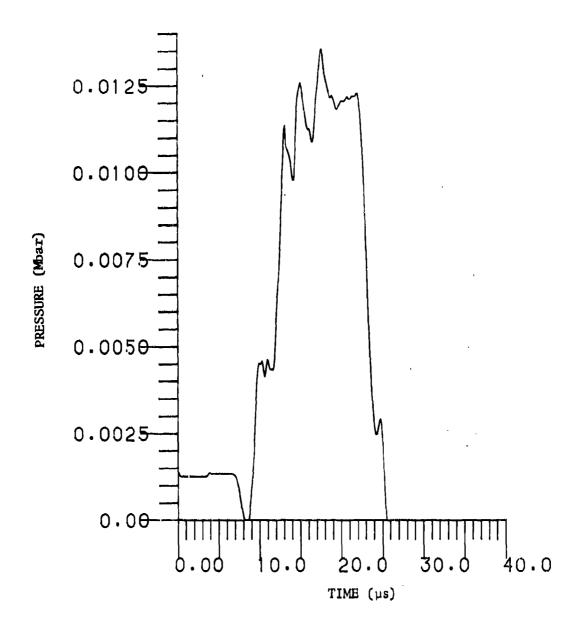


Figure 10c. Acceptor Explosive Pressure History with 30-mm Thick
Three-Layered Tungsten/Lucite/Tungsten Shield at 33.3% Lucite

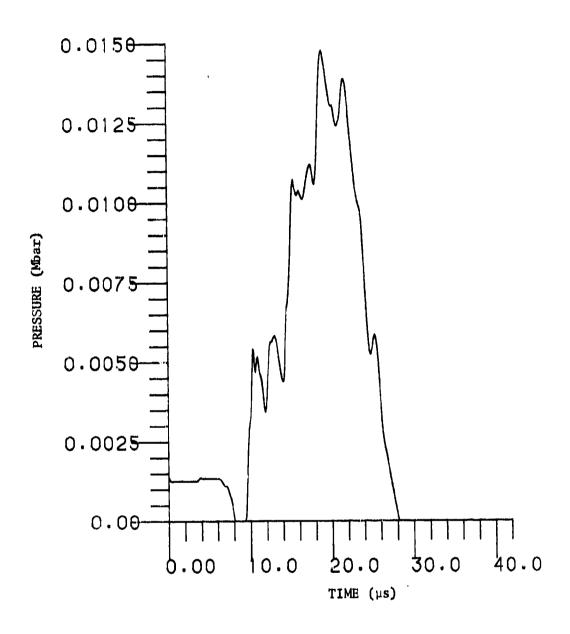


Figure 10d. Acceptor Explosive Pressure History with 30-mm Thick Three-Layered Tungsten/Lucite/Tungsten Shield at 50% Lucite

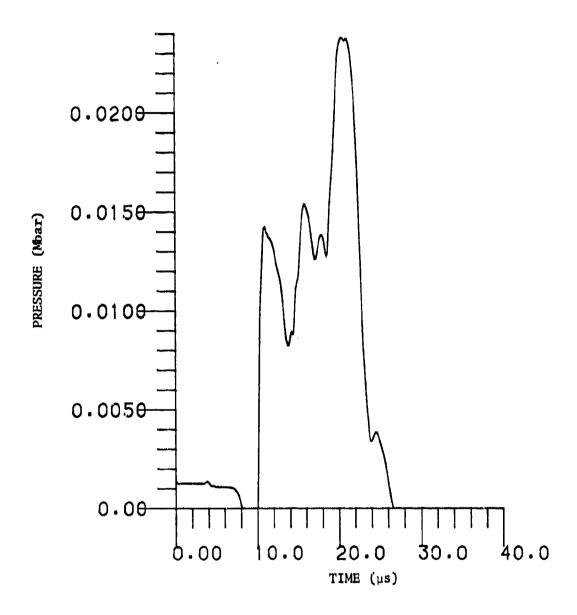


Figure 10e. Acceptor Explosive Pressure History with 30-mm Thick
Three-Layered Tungsten/Lucite/Tungsten Shield at 91.7% Lucite

may be noted but this is still primarily a single shock. With a still higher Lucite fraction, the shock structure is altered radically and multiple shock loading is observed as shown in Figure 10c through 10e. The transition to multiple shocks occurs somewhere between 17 and 33 percent Lucite and minimum pressure levels occur at about 33 percent Lucite.

When tungsten is the inner component of the shield, the results are somewhat different as illustrated in Figure 11. In this case a single pulse is usually produced but the peak pressure and rise time of the compression depend upon the Lucite thickness fraction as indicated in Table 2. This shows the pressurization rate to be minimum at about 50 percent Lucite. Multiple shock structure does not appear until the shield is nearly all Lucite as shown in Figure 11d.

Table 2. Compression Wave Characteristics for Lucite/Tungsten/Lucite Shields

Shield Configuration (mm/mm/mm)	Lucite Fraction	Peak Pressure (GPa)	Rise Time (µs)	Pressurization Rate (GPa/µs)
1.25/27.5/1.25	.083	3.95	1.6	2.47
2.5/25.0/2.5	.167	3.28	3.5	0.94
5.0/20.0/5.0	.333	2.57	4.1	0.42
7.5/15.0/7.5	.500	2.40	10.0	0.24
10.0/10.0/10.0	.667	3.06	11.0	0.28
2.5/ 5.0/12.5	.833	3.81	9.5	0.40
13.75/ 2.5/13.75	.917	3.89	5.0	0.78

The shock structure observed with steel/Lucite shields is essentially the same as that for the tungsten/Lucite shields with somewhat higher pressure levels. Therefore, pressure history plots are not reproduced here. The effect on compression wave characteristics for Lucite/steel/Lucite shields is summarized in Table 3. The pressurization rate is again minimized for about 50 percent Lucite.

Table 3. Compression Wave Characteristics for Lucite/Steel/Lucite Shields

Shield Configuration (mm/mm/mm)	Lucite Fraction	Peak Pressure (GPa)	Rise Time (µs)	Pressurization Rate (GPa/µs)
1.25/27.5/1.25	.083	5.18	1.4	3.70
2.5/25.0/2.5	.167	4.61	2.1	2.20
5.0/20.0/5.0	.333	3.62	5.9	0.61
6.3/17.4/6.3	.420	3.62	7.0	0.52
7.5/15.0/7.5	.500	3.92	8.4	0.47
10.0/10.0/10.0	.667	4.10	6.2	0.67
12.5/ 5.0/12.5	.833	3.90	5.8	0.67
13.75/ 2.5/13.75	.917	3.43	5.0	0.69

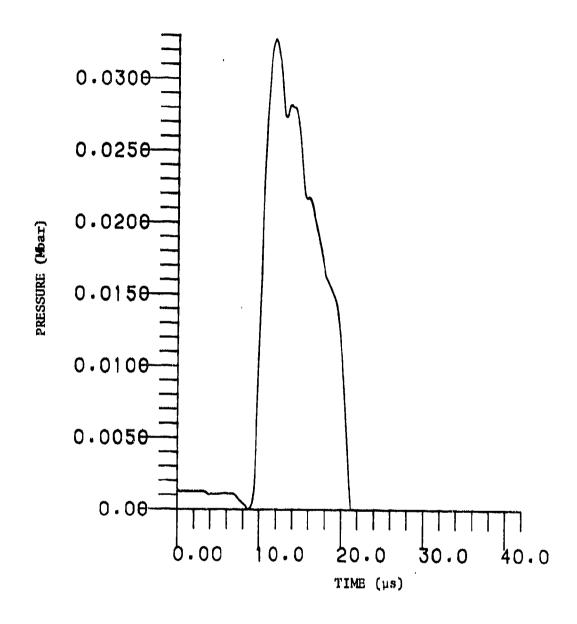


Figure 11a. Acceptor Explosive Pressure History with 30-mm Thick Three-Layered Lucite/Tungsten/Lucite Shield at 16.7% Lucite

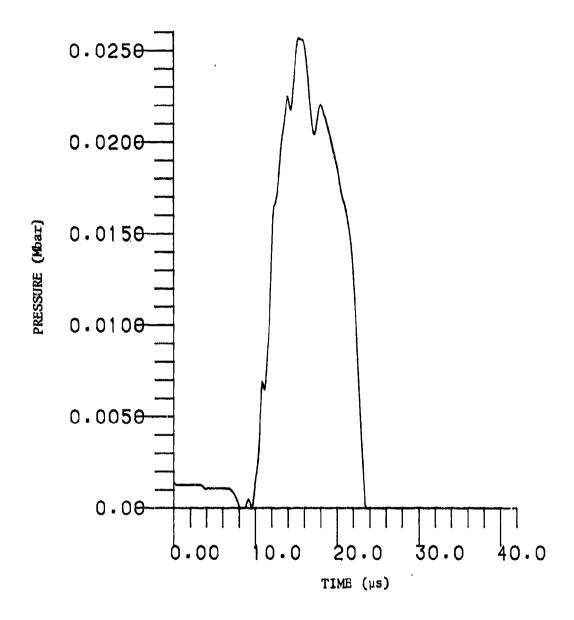


Figure 11b. Acceptor Explosive Pressure History with 30-mm Thick Three-Layered Lucite/Tungsten/Lucite Shield at 33.3% Lucite

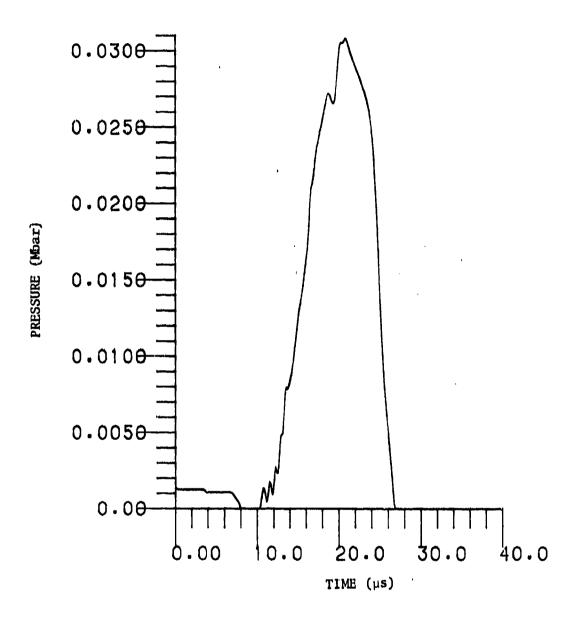


Figure 11c. Acceptor Explosive Pressure History with 30-mm Thick
Three-Layered Lucite/Tungsten/Lucite Shield at 66.7% Lucite

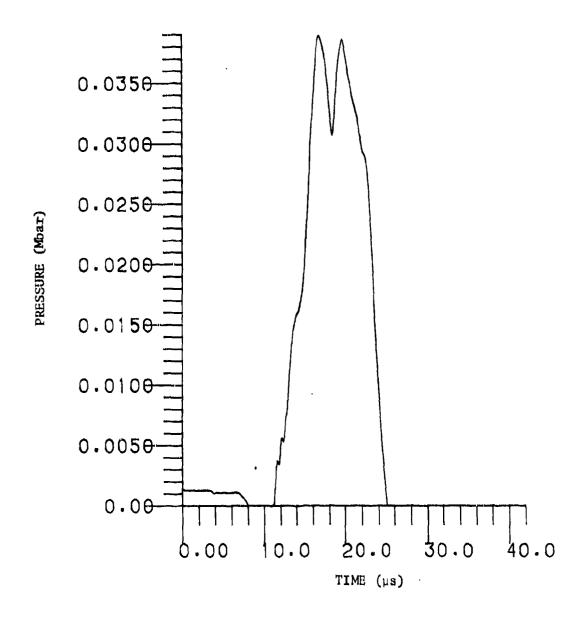


Figure 11d. Acceptor Explosive Pressure History with 30-mm Thick Three-Layered Lucite/Tungsten/Lucite Shield at 91.7% Lucite

Increasing shield thickness has considerable effect on shock breakup. This is illustrated in Figure 12 for 15, 25, 35, and 45 mm thick three-layered shields having tungsten as the outer component. Thin shields produce no breakup, presumably because the close spacing between impedance discontimuities produces closer temporal spacing between shocks leading to rapid coalescence. Figure 12a, for a 15 mm thick shield, seems to show shocks which have just coalesced resulting in a compression wave with a higher rise time. On the other hand, Figure 12d, for a 45 mm thick shield, shows the most shock breakup and peak pressure reduction we have seen in any of our computations.

Five-layered shields produced considerably less shock breakup than three-layered shields of the same thickness, both having tungsten as the outer component. This can be seen by comparing Figure 10d, for a 50 percent Lucite three-layered shield, with Figure 13, for a 50 percent Lucite five-layered shield. Again this presumably occurs because of the closer spacing between impedance discontinuities. This five-layered shield does produce a rise time of 6 to 7 ms.

C. <u>Initiation Stimulus Reduction</u>

Consideration of stimulus levels as characterized by $\int p^2 dt$ allows a comparison of the effectiveness of these shields. Figure 14 is a plot of the relative stimulus level ($\int_{p}^{2} dt$ normalized with respect to the value for the unshielded case) versus Lucite thickness fraction for tungsten/Lucite/tungsten and Lucite/tungsten/Lucite shields. Also included are scales of Lucite weight fraction and areal density. This comparison also shows that considerably better shield performance is obtained when tungsten is the outer component of the "sandwich." In fact, the performance of Lucite/tungsten/Lucite shields is inferior to that of either pure tungsten or pure Lucite except for a range of Lucite thickness fraction values between about 0.2 and 0.6 wherein performance is only a little better. It should be noted that this comparison is based on $\int_{\rm P}^2$ dt and does not take into account the differences between the ramp compressions produced when Lucite is the outer component and the shock breakup and lower peak pressure produced when tungsten is the outer component. Indeed, the response of explosives to these complex waveforms is not well understood and experimental verification is required to determine which shield configuration is superior. However, it should be remarked that the predicted advantage of shields with tungsten on the outside is substantial.

Similar observations may be made for three-layered steel/Lucite shields as shown in Figure 15. The performance of tungsten/Lucite/tungsten shields is compared with that of steel/Lucite/steel shields in Figure 16. This shows that shields incorporating tungsten reduce the initiation stimulus to levels well below those produced by shields incorporating steel.

Shield thickness was varied in two different ways. Shields were varied from 10 to 50 mm in thickness with a constant 50 percent Lucite fraction in one case and from 20 to 50 mm with a constant 7.5 mm outer tungsten layer thickness in the other. The results are plotted versus shield thickness in Figure 17. This shows little difference between the two types of variation and indicates a rapid decline in protection for shields less than about 30 mm thick.

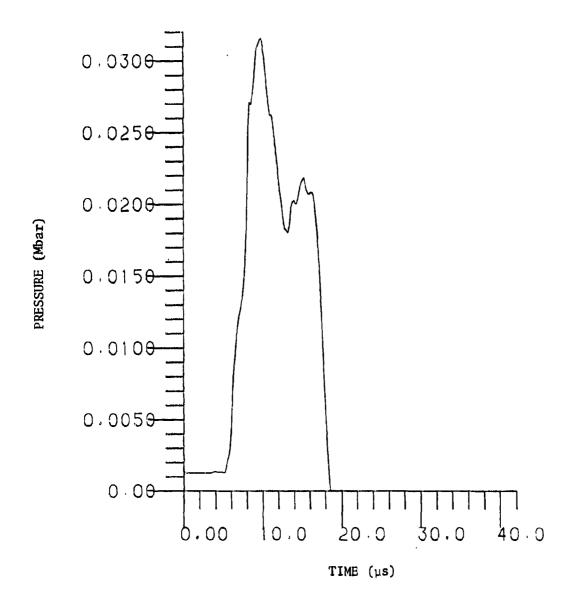


Figure 12a. Acceptor Explosive Pressure History with 50% Lucite Three-Layered Tungsten/Lucite/Tungsten Shields at 15-mm Shield Thickness

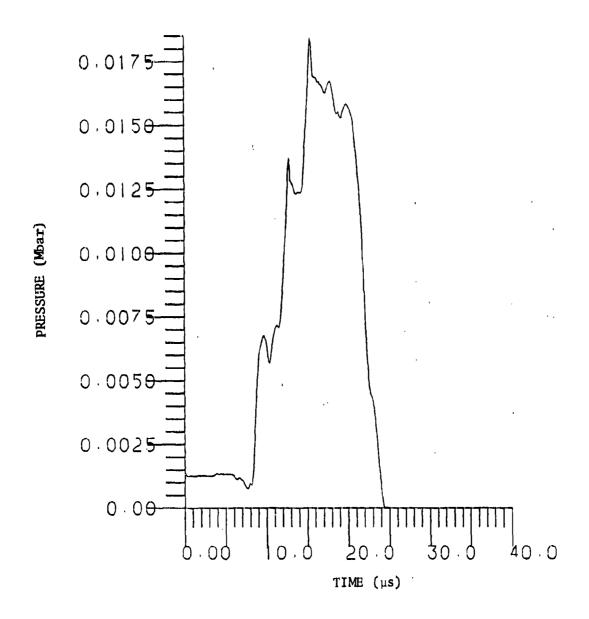


Figure 12b. Acceptor Explosive Pressure History with 50% Lucite Three-Layered Tungsten/Lucite/Tungsten Shields at 25-mm Shield Thickness

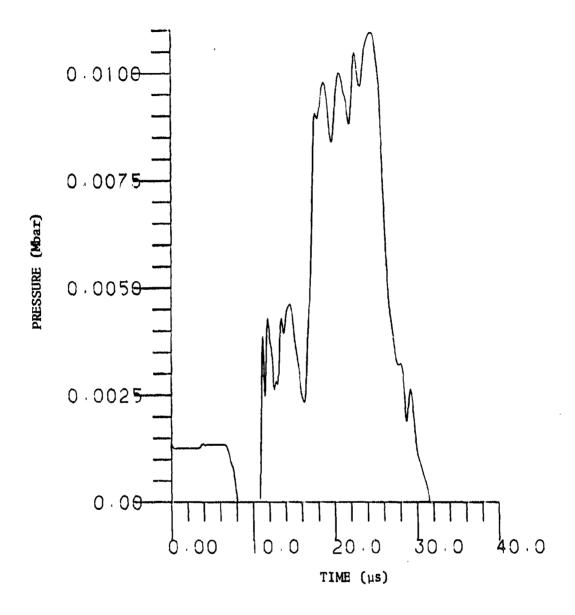


Figure 12c. Acceptor Explosive Pressure History with 50% Lucite Three-Layered Tungsten/Lucite/Tungsten Shields at 35-mm Shield Thickness

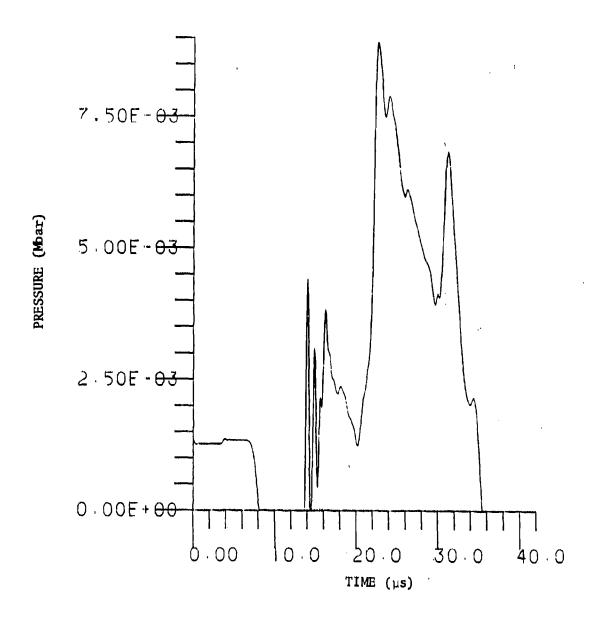


Figure 12d. Acceptor Explosive Pressure History with 50% Lucite Three-Layered Tungsten/Lucite/Tungsten Shields at 45-mm Shield Thickness

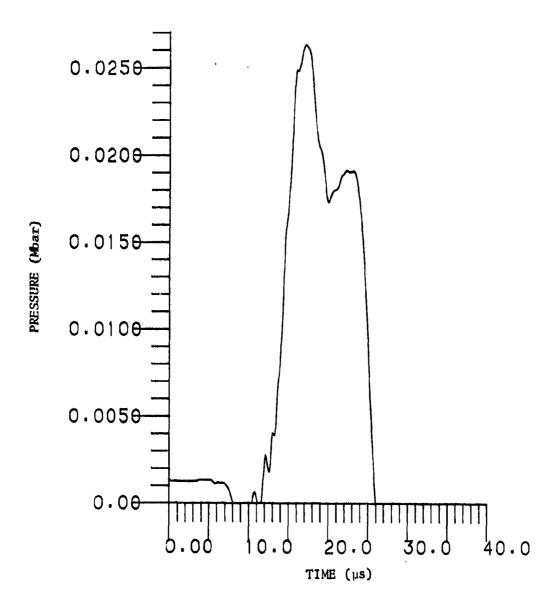


Figure 13. Acceptor Explosive Pressure History with a 30-mm Thick, 50% Lucite Five-Layered Tungsten/Lucite/Tungsten/Lucite/Tungsten Shield

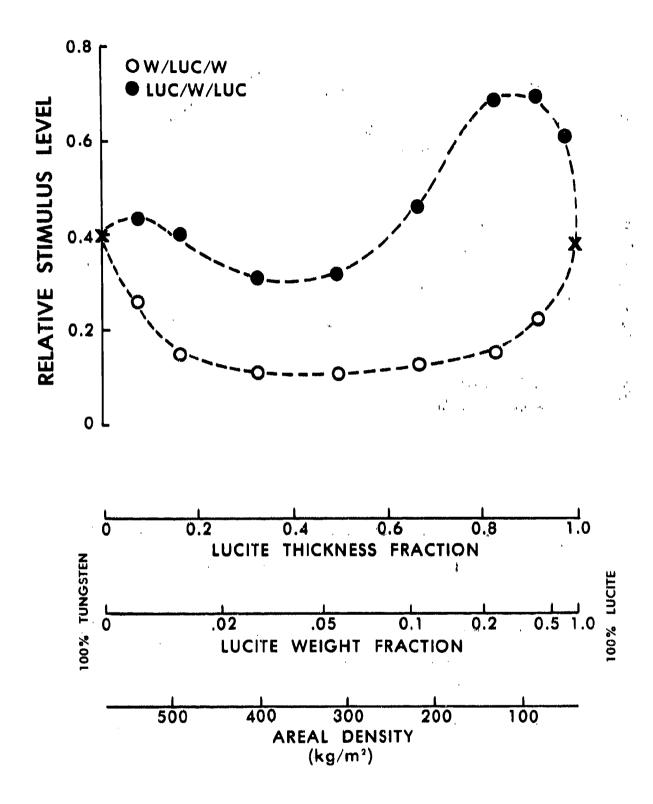


Figure 14. Relative Stimulus Levels for 30-mm Thick, Three-Layered Tungsten/Lucite Shields

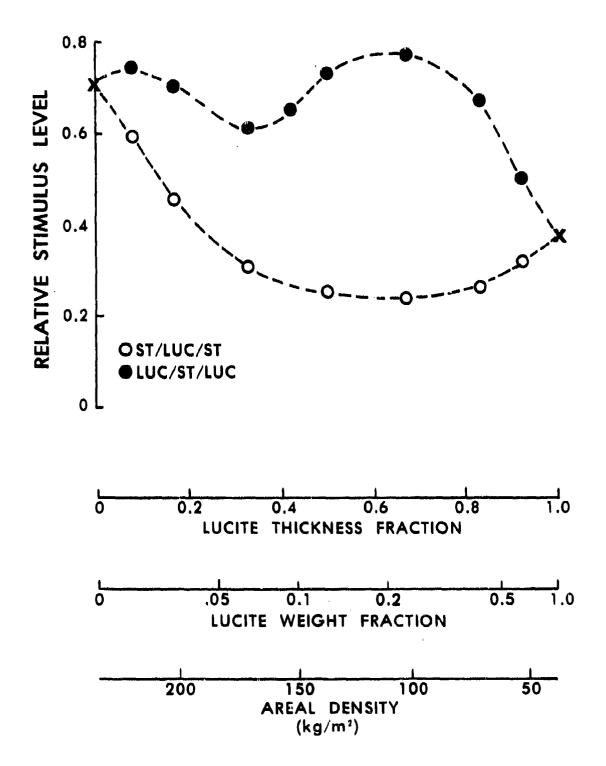
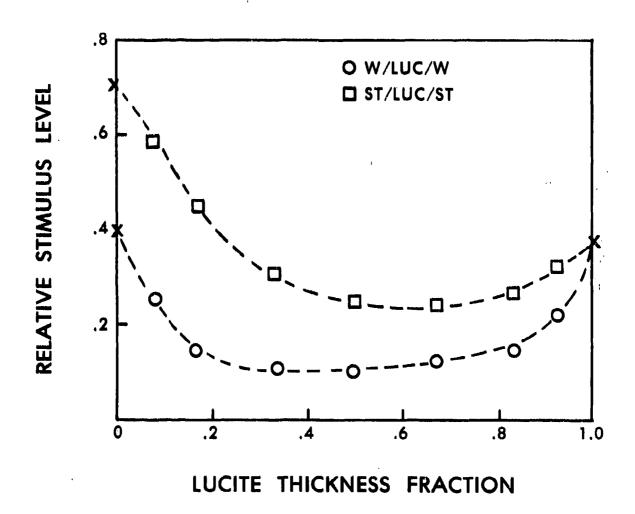


Figure 15. Relative Stimulus Levels for 30-mm Thick, Three-Layered Steel/Lucite Shields



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Figure 16. Comparison of Relative Stimulus Levels for 30-mm Thick, Three-Layered Tungsten/Lucite and Steel/Lucite Shields

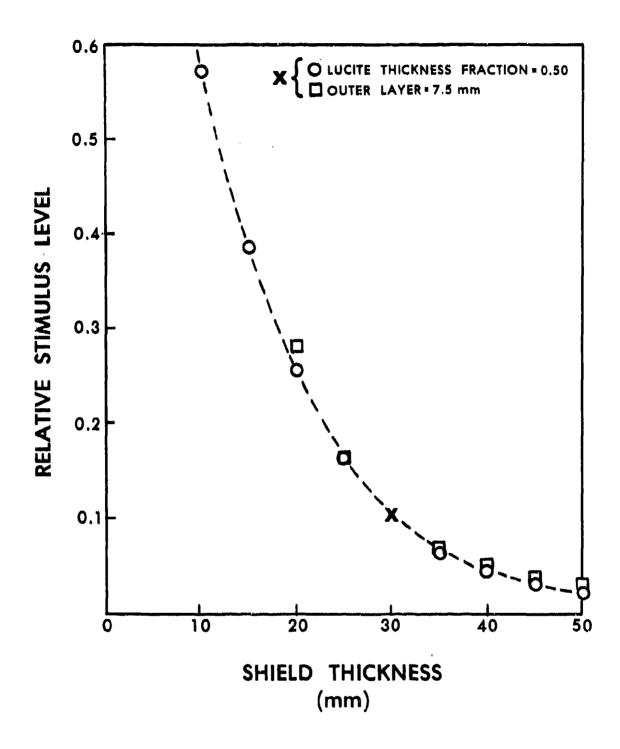


Figure 17. Effect of Shield Thickness on Relative Stimulus Level for Three-Layered Tungsten/Lucite/Tungsten Shields

Results with five-layered shields were similar to those for three-layered shields but not as much protection was provided. We first considered shields with inner and outer layers of equal thickness (h,=h,). Figure 18 shows the comparison between shields with tungsten on the outside and shields with Lucite on the outside. Again, the former perform better but only slightly better than single layered shields. Figure 19 shows the comparison between the best three-layered shields and the five-layered shields. The threelayered shields are substantially better. The performance of the five-layered shields with tungsten as the outer component is not strongly dependent on Lucite fraction. It is of interest to determine whether shields with unequal inner and outer layer thicknesses perform any better. We noted previously that the parameter z could be used as a measure of the relationship between these layers and that z=0 and z=h4/h correspond to three-layered shields. We fixed the Lucite thickness fraction at 0.33 and varied z between these limits. The results, shown in Figure 20, indicate a nonmonotonic variation between the two three-layered cases. The three-layered shield with tungsten on the outside still shows the best performance.

VII. SUMMARY

We have conducted a study of the role of shielding in reducing the shock initiation stimulus for a simple one-dimensional representation of the problem of sympathetic detonation of munitions. We found that single layered shields made of materials with low acoustic impedance generally produce a complex shock wave structure in the acceptor. This complex structure is associated with a low level of initiation stimulus because of the breakup into several weaker shocks and the reduction of $\int p^2 dt$. High impedance shield materials also substantially reduce the initiation stimulus, but without the accompanying shook breakup effect. Increasing shield thickness improves performance and can change the order of effectiveness of shield materials. With multilayered shields composed of a high-impedance and a low-impedance material we observed shock structures depending strongly on the ordering of the materials in the shield. Multiple shock structure was usually observed when the highimpedance material was the outer component of the shield and sufficient lowimpedance material was present. Single compression waves with variable peak pressures and rise times were usually observed when the low-impedance material was the outer component of the shield. Substantial benefits in terms of shock breakup and $\int p^2 dt$ reduction can be obtained by increasing the thickness of three-layered shields, which were found to perform better than five-layered shields.

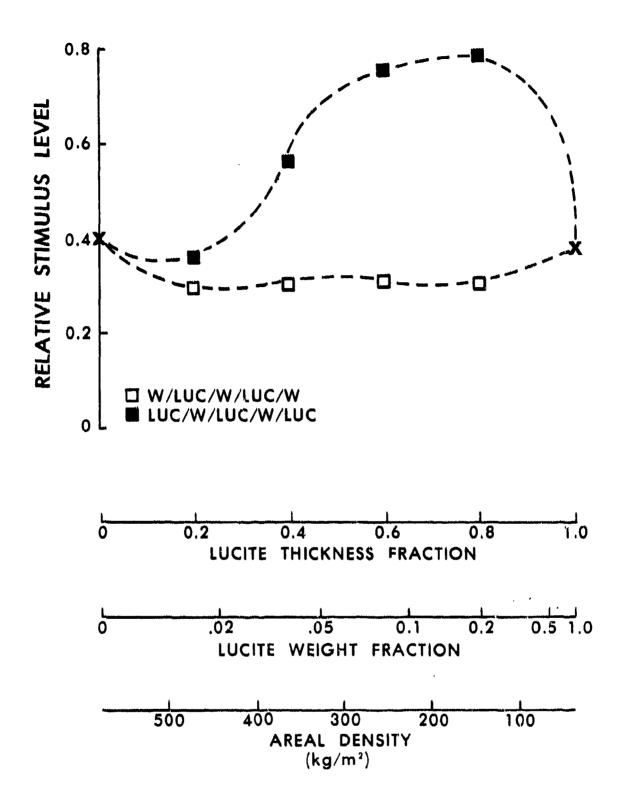
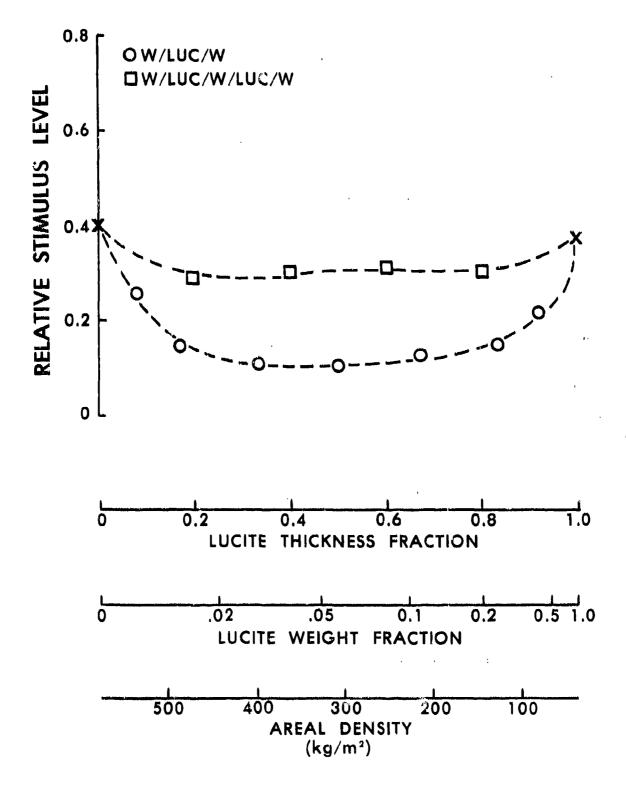


Figure 18. Relative Stimulus Levels for 30-mm Thick, Five-Layered Tungsten/Lucite Shields



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Figure 19. Comparison of Relative Stimulus Levels for 30-mm Thick, Three-Layered and Five-Layered, Tungsten/Lucite Shields

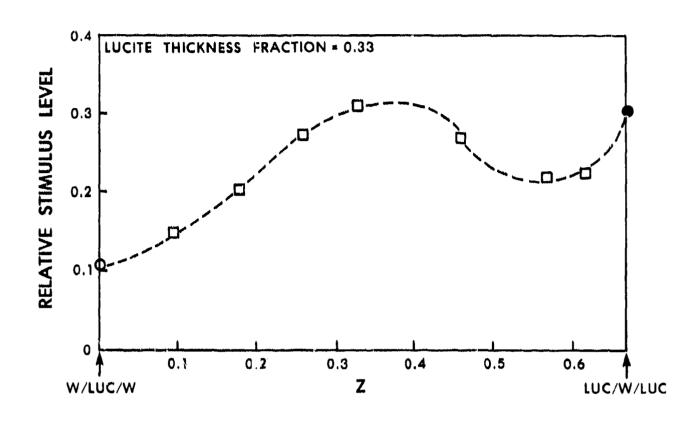


Figure 20. Effect of Shield Configuration on Relative Stimulus Level for 33.3% Lucite, Three-Layered and Five-Layered Shields

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